Changes in Cloud Optical Thickness and Cloud Drop Size Associated with Precipitation Measured with TRMM Satellite

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Abstract

We derived rain rate, cloud optical thickness and the effective radius in water clouds by a combined use of the Precipitation Radar and the Visible and Infrared Scanner onboard the Tropical Rainfall Measuring Mission. The derived data were used to study how cloud optical thickness relates to precipitation. In particular, we focused on the changes in cloud optical thickness resulted from changes in the size distributions of cloud droplets associated with precipitation. There were considerable scatter between cloud optical thickness and rain rate on a global scale. However, cloud optical thickness was found to increase with rain rate on average. The tendency to increase was mostly due to increases in liquid water path and depended on rain rate. For strong rain, relatively small increases in the optical thickness with rain rate were observed. Whereas, for weak rain, larger increases with rain rate were found, which is related to considerable changes in liquid water path and in the effective radius of cloud droplets. To study the effects of drop size variation, the relationships between cloud optical thickness and rain rate for same values of liquid water path were analyzed. Results show that there were no significant dependences of cloud optical thickness on rain rate for strong rain. For weak rain, cloud optical thickness was found to decrease with rain rate. In particular, significant differences of optical thickness were found between non-precipitating clouds and precipitating clouds: smaller cloud optical thickness was observed for precipitating clouds. Dispersion of cloud drop size was found in the rain formation process, which may relate to changes in the shape of drop size spectra and leads to the decreases in the cloud optical thickness for precipitating clouds.

1. Introduction

The cloud feedback problem is one of the largest uncertainties in climate studies (Stephens 2005) because cloud formation and dissipation are complicated processes. Clouds have mutual relations to aerosols and precipitation. Increases in aerosol concentrations result in a decreased drop size and modify the cloud radiative forcing, which is known as the 'first indirect effect' (Twomey 1974; Breon et al. 2002). The interactions between aerosols and clouds have been studied extensively using numerical models and observations with airborne sensors (e.g., Kuba et al. 2003; Asano et al. 1995) and space-borne solar and infrared radiometers (Kobayashi and Masuda 2008; Coakley and Bernstein 1987). Decreased drop size may also suppress precipitation, which is known as the 'second indirect effect' (Albrecht 1989; Rosenfeld 2000). Precipitation is generated from cloud drops by condensation and coalescence, which in turn generates cloud drops by breakup (Kobayashi and Adachi 2001), and removes cloud drops and controls cloud amount. A study on clouds-precipitation interaction is, therefore, critical to improve our understanding of the cloud feedback problem. Despite extensive studies on cloud-precipitation interactions (e.g., vanZanten et al. 2005), our understanding is very limited because of their complex nature.

A recent study suggested that there exist significant differences of the effective radius \(r_e\) of cloud drops between non-precipitating clouds (NPC) and precipitating clouds (PRC) (Kobayashi 2007). The
showed that increases in \( t \) arises from increased \( \tau \) (Gerber 1996), which compensates the decreases in the critical radius drops grow to raindrops quickly when they exceed differences are associated with the fact that cloud-precipitation interactions. The orbit has an inclination of 35° and the observed area is from \(-35°\) to \(+35°\) in latitude and \(0°\) to \(360°\) in longitude. We selected the VIRS data for which the center position of the PR is within the footprint of the VIRS and made match-up datasets from data for 11 days in December, 2002 and 10 days in June, 2003. The swath of the VIRS is wider than that of the PR. The VIRS data extending beyond the swath of the PR were excluded. Ultimately, we generated two global-segmented datasets, each with data for 10 or 11 days. One segment box has a \(0.1° \times 0.1°\) latitude-longitude spatial resolution. Data exist for about 4~5 pixels in each grid box. The pixel for precipitating at the highest rain rate or for the highest reflection at Ch1 (0.63 \(\mu m\)) when no precipitating pixels occur in a box was selected. Horizontally inhomogeneous clouds and cirrus clouds lead to some error (Kobayashi 1988; Kobayashi 1993; Han et al. 1994). In the analysis, pixels for which the brightness temperature at Ch4 (10.8 \(\mu m\)) of VIRS ranged from 273 to 290 K and the brightness temperature difference between Ch4 and Ch5 (12 \(\mu m\)) was less than 1 K were furthermore selected to remove erroneous data of optically thin, cold clouds and sub-pixel clouds (Rosenfeld and Gutman 1994).

3. Method

Cloud optical thickness and the effective radius can be estimated from space-borne measurements of radiances at two wavelengths based on the differences in water absorbing characteristics, like 0.75 \(\mu m\) and 2.16 \(\mu m\) (Nakajima and King 1990; Han et al. 1994) or more accurately by 0.64 \(\mu m\) and 3.75 \(\mu m\) (Nakajima and Nakajima 1995). Here, the reflected radiances at Ch1 (0.63 \(\mu m\)) and Ch3 (3.75 \(\mu m\)) measured with the VIRS were used to de-
rive the optical thickness and the effective radius of cloud droplets near the cloud tops (Kobayashi 2007). The effective radius is defined as follows.

\[
Re = \frac{\int r^2 N(r) \, dr}{\int r^2 \, dr}
\]  

(1)

where, \(N(r)\) is drop size distribution and \(r\) is radius of droplets. The retrieval method is based on the difference of water absorbing characteristics between solar reflectance at Ch1 and Ch3. The size distribution of cloud droplets is assumed to be the lognormal function as:

\[
N(r) = \frac{N_0}{\sqrt{2\pi\sigma}} \exp\left[-\frac{(\log r - \log r_0)^2}{2\sigma^2}\right].
\]  

(2)

Here, \(r_0\) is the mode radius and the variance \(\sigma\) is assumed to be 0.35 (Nakajima and King 1990). The phase function is calculated by Mie theory assuming spherical shape of water droplets. Bidirectional reflections at the top of atmosphere are calculated for various values of \(r_e\) and \(\tau\) by using radiative transfer model.

The radiance at Ch3 includes contributions of thermal emissions from the atmosphere and ground surface. The undesired thermal emissions from the cloud tops were removed from the measured radiance at Ch3. The cloud top height was determined from radiance at Ch4 (10.8 \(\mu\)m) using LOWTRAN-7 (Kneizys 1988) assuming the tropical model for the atmosphere above the cloud top. This assumption leads to some error in the derived \(r_e\). However, a difference of 4 K between the model and the actual temperature profile causes an error of less than 0.7 \(\mu\)m in the retrieved \(r_e\) (Han et al. 1994; Kawamoto et al. 2001). Use of the mid-latitude summer model for the atmosphere above the cloud top instead of the tropical model results in a difference of 0.8 \(\mu\)m in the mean \(r_e\). Since we applied the method to limited regions of latitude, this error is not serious. We ignored the thermal emission from the ground surface because we examined optically thick clouds (King 1987).

4. Results

Figure 1 shows the relationships between mean cloud optical thickness and rain rate for data from \(0.1^\circ \times 0.1^\circ\) latitude/longitude grids (dotted line). Data from \(0.5^\circ \times 0.5^\circ\) grids are also plotted for comparison (solid line). Although there are large standard deviations of \(\tau\) (vertical bars, upper: \(0.1^\circ \times 0.1^\circ\), lower: \(0.5^\circ \times 0.5^\circ\)), cloud optical thickness clearly tends to increase with rain rate for weak rain and tends to be constant and saturated for strong precipitation. Large error bars may interrupt reader’s understanding, therefore we omitted error bars for other figures. Sharp increases in \(\tau\) are observed between NPC (rain rate = 0) and PRC for a segment box of \(0.1^\circ \times 0.1^\circ\) which contains only 3–5 pixels. For a segment box of \(0.5^\circ \times 0.5^\circ\), however, cloud optical thickness is relatively large for NPC. Cloud optical thickness decreases in the conversion process from NPC to PRC with rain rate of 0.4 mm/h. Data exist for at least 100 pixels in each grid box of \(0.5^\circ \times 0.5^\circ\). When no precipitating pixels occur in a box, the highest reflection at Ch1 was selected as a grid data. Consequently, optically thicker clouds were selected for \(0.5^\circ \times 0.5^\circ\) grid box than those for \(0.1^\circ \times 0.1^\circ\) grid. The decreases in \(\tau\) found for \(0.5^\circ \times 0.5^\circ\) grid data may be, therefore, artifact. However, measured non-precipitating clouds from \(0.5^\circ \times 0.5^\circ\) grid box more likely represent NPC that grows to PRC than those from \(0.1^\circ \times 0.1^\circ\) grid box, because more vigorous clouds produce precipitation. Thus, the decrease in \(\tau\) may really occur in the conversion process from cloud droplets to raindrops. In the onset of precipitation, number of cloud droplets likely decreases in the conversion process associated with rapid growth by collision and removal of cloud droplets by drizzle. The de-
creased droplets concentration results in the decreases in cloud optical thickness for PRC. These effects are critical to examine the second indirect effect in which increased aerosols result in suppression of precipitation and affect the lifetime of clouds.

Cloud optical thickness is determined primarily by liquid water path (LWP) and secondarily by drop size distributions as mentioned earlier. Size of cloud droplets changes by factor of 1 or 2, whereas LWP ranges over factors of 2 or 3 and primarily determines \( \tau \). Liquid water path can be derived as a secondarily variable from primary observables, \( \tau \) and \( r_e \) approximated as

\[
\text{LWP} \approx \frac{2}{3} \tau \rho \text{Re},
\]

where, \( \rho \) is the density of liquid water. Figure 2 shows variations of \( \tau \) and LWP with rain rate. Relative values of LWP and \( \tau \) to those at rain rate = 0 are plotted. The relative values of LWP increase more significantly than those of \( \tau \) do. In particular for rain rate ranging from 0 to 1 mm/h, increases in \( \tau \) are slower than LWP. Cloud optical thickness is linearly related to LWP for a fixed drop size distribution. Thus, the different increasing tendencies between \( \tau \) and LWP suggest that there exist changes in \( r_e \) and likely in the spectral shape of the size distributions of cloud droplets with rain rate. In particular, large differences are found for degree to which \( \tau \) and LWP increase with rain rate in the conversion process from cloud droplets (NPC) to raindrops (PRC). This indicates that there are significant differences of the size distributions that are active in the visible optics region between NPC and PRC.

In the present study, we intend to examine the effects of precipitation on \( \tau \) in terms of the size of cloud droplets. We, therefore, need to remove the contribution of LWP to \( \tau \). To study the effects of drop size on \( \tau \), we examined the relationships between \( \tau \) and rain rate for fixed values of LWP. We averaged \( \tau \) for three values of LWP ranging from 0 to 130, from 130 to 260 and from 260 to 400 g/m² (Fig. 3). Cloud optical thickness tends to decrease with rain rate, which is in contrast to the increasing tendencies of \( \tau \) with rain rate as shown in Fig. 1. The decreasing tendencies of \( \tau \) with rain rate for fixed values of LWP shown in Fig. 3 are mostly due to the change in the size distributions of cloud droplets. For stronger rain, cloud optical thickness is almost constant as in Fig. 1, suggesting that no significant changes in the corresponding size distributions occur within the size range detected by visible optics.

Figure 4 shows the dependence of \( r_e \) on rain rate for three ranges of LWP as in Fig. 3. The effective radius is almost constant for rain rate larger than 1 mm/h, corresponding to the small variations of the ratio of \( \tau \) to LWP with rain rate as shown in Fig. 3. For weak rain, \( r_e \) increases with rain rate, corresponding to the decrease in \( \tau \) with rain rate as
shown in Fig. 3. Previous study reported significant increases in $r_e$ for PRC comparing with NPC (Kobayashi 2007). Figure 4 shows that the significant differences in $r_e$ between NPC and PRC occur for any values of LWP.

The increases in $r_e$ with rain rate likely occur associated with changes in the shape of the size distribution of cloud droplets. For NPC, most cloud droplets are smaller than $r_e$ and primarily grow by condensation. Whereas, for PRC, cloud droplets larger than $r_e$ primarily grow by coalescence. These different processes of growth possibly result in changes in the cloud drop size distributions between NPC and PRC. During drizzle formation, broadening of the drop spectra likely occurs (Feingold et al. 1997). The spectral broadening to larger size leads to increase in $r_e$ and decreases in $\tau$ assuming constant LWP. The spectral broadening, therefore, is a key factor to improve our understanding of cloud-precipitation interaction. It is needed to examine whether the spectral broadening occurs on a global scale or not. In the present study, however, we derived $\tau$ and $r_e$ by assuming the lognormal function with fixed value of $\sigma$ as 0.35 for the cloud drop spectra and cannot determine the variations of the spectral broadening explicitly. To estimate the dispersion of the cloud droplets size distribution, we have used the variations of the derived values of $r_e$ instead. The variation of the derived effective radius derived for pixel by pixel can be a measure of the dispersion of the averaged size distribution for a large area to some degree.

Figure 5 shows the variances of $r_e$ as a function of rain rate for various values of LWP as in Fig. 3. Significant increases in the mean variances (dashed line) over whole values of LWP are found at weak rain rate. For large values of LWP, the increases are also found at weak rain rate. Whereas, no apparent peaks are found for small LWP. These features well correspond to the variation of $\tau$ with rain rate shown in Fig. 3. Vertical variations of the standard deviation of cloud droplet mean diameter are larger for higher altitude in which more drizzle and broadened drop size distribution are observed (Hudson and Yum 2001). The increases in the variance of $r_e$ shown in Fig. 5, are therefore thought to be associated with broadened size distribution.

Figure 6 shows average drop size distributions obtained from the derived effective radius over rain rate ranging from 0 to 0.3, 0.5 to 1, and 3 to 4 mm/h as well as NPC. Normalized size distributions are plotted. The width of the size distribution is narrower for NPC than that for PRC. For PRC, the mode radius is larger and the width of the distributions is wider for stronger precipitation. For rain rate larger than a few mm/h, no significant changes in the size distributions appear, suggesting increases in large cloud droplets that are inactive in the visible optics region associated with strong precipitation. For rain rate larger than a few mm/h, no significant changes in the size distributions appear, suggesting increases in large cloud droplets that are inactive in the visible optics region associated with strong precipitation. Differences of the normalized size distributions between NPC (solid lines) and PRC (dotted lines) are shown in Fig. 7 for LWP ranging from 0 to 130 and 260 to 400 g/m$^2$. For small value of LWP, slight differences are found. For larger LWP, however, broader shape of drop size distribution is
found for PRC, corresponding to the significant change in the variations shown in Fig. 5. Figure 8 shows relative values of cloud optical thickness calculated theoretically by using Mie theory from the averaged drop size distributions as a function of rain rate for various values of LWP. Observed optical thickness (Fig. 3) relative to that at rain rate = 0 are also plotted (circles, crosses and triangles). Relative values of \( \tau \) were calculated from the normalized size distributions so that they were for the same liquid water content. Decreases in \( \tau \) are clearly found in clouds with weak rain, which is in good agreement with the observations. Figure 8 suggests that the variations of cloud optical thickness with rain rate are closely related to changes in the shape of size distributions of cloud droplets.

5. Conclusions

Cloud optical thickness and the effective radius were derived from a combined use of the VIRS and the PR onboard the TRMM satellite to examine the effects of precipitation on the cloud properties. We focused on the effects of the size distributions of cloud droplets on cloud optical thickness associated with precipitation. Overall cloud optical thickness increases with rain rate for weak rain and is almost independent of rain rate for stronger rain than a few mm/h. The increase in cloud optical thickness was found to be primarily due to increases in LWP. To exclude the effects of increases in LWP on \( \tau \), cloud optical thickness for fixed values of liquid water path was examined. Such derived cloud optical thickness for weak rain showed opposite dependence of rain rate for large values of LWP. Cloud optical thickness tended to decrease with rain rate and was larger for NPC than that for PRC. In particular for large values of LWP, difference of \( \tau \) between NPC and PRC was
significant. Dispersion of cloud drop size was found in the rain formation process, which may relate to changes in the shape of drop size spectra and leads to the decreases in the cloud optical thickness for precipitating clouds.

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