Voronoi Analysis of the Structure of Ni-Zr-Al Ternary Metallic Glass

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Ni-Zr metallic glasses have been recognized to be unstable in comparison with Cu-Zr metallic glasses. An analysis of Voronoi polyhedra in the RMC simulations based on the diffraction data could characterize the atomic configurations around Ni and Cu atoms. The polyhedra around Ni atoms are dominated by trigonal prism-like, Archimedian antiprism-like, and similar polyhedra. In contrast, icosahedron-like polyhedra are preferred for Cu. The Ni-Zr glasses have been reported to stabilize by adding Al. Therefore, in this work, the analysis of Voronoi polyhedra around Ni, Zr and Al atoms for Ni₂₅Zr₆₀Al₁₅ ternary metallic glass was carried out in order to clarify the difference between the atomic structures for the binary and ternary metallic glasses. Trigonal prism-like, Archimedian antiprism-like and similar polyhedra, which are dominated in the Ni-Zr metallic glasses, decreased in number by adding Al to the Ni-Zr system. On the contrary, the number of icosahedron-like polyhedra was found to increase. The results apparently indicate that the addition of Al into Ni-Zr binary system promote the formation of icosahedron-like polyhedra in the structure. Therefore, from these results, we can easily recognize that icosahedron-like polyhedra play an important role to stabilize the structure of metallic glasses. [doi:10.2320/matertrans.MJ200750]

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1. Introduction

Bulk metallic glasses have received much attention in order to have a large number of technical applications as high strength structural materials. It is well known that the metallic glasses become stable by addition of third, forth and fifth elements to the binary system. For example, it has been reported that the glass transition in the Ni-Zr metallic glass is observed in a wide compositional range by addition of small amount of Al.¹⁾ Structural investigation for Ni-Zr-Al metallic glasses has informed us of the change in the local structure by addition of Al.²⁾ That is, the coordination number of 12.5 around Zr for Zr₆₀Ni₂₅Al₁₅ metallic glass is much different from that of 10.7 for Zr₇₀Ni₃₀ metallic glass. However, the result only suggests the formation of icosahedral local configuration in Zr₆₀Ni₂₅Al₁₅ metallic glass and no detailed topological information has reported.

In this work, we have investigated the topology of the atomic configurations in Ni_{33.3}Zr_{66.7} binary and Ni₂₅Zr₆₀Al₁₅ ternary metallic glasses by using a combination of neutron & X-ray diffraction data and Reverse Monte Carlo (RMC) modeling.³⁾ Time-of-flight (T-O-F) neutron diffraction using neutrons of short-wavelength and X-ray diffraction with photon energy of 113.68 keV were employed to make high real-space resolution measurements on the atomic configurations of Ni_{33.3}Zr_{66.7} and Ni₂₅Zr₆₀Al₁₅ metallic glasses. RMC modeling has been recognized to be an excellent method for visualizing the three-dimensional atomic arrangement of liquid and amorphous materials, based on the results of neutron & X-ray diffraction experiments. Therefore, in this study RMC modeling based on the neutron & X-ray diffraction data was carried out for constructing the atomic configurations of the Ni_{33.3}Zr_{66.7} and Ni₂₅Zr₆₀Al₁₅ metallic glasses and for investigating the topological characteristics.

2. Experimental

Ni_{33.3}Zr_{66.7} binary and Ni₂₅Zr₆₀Al₁₅ ternary alloys were prepared by arc melting using Ni (99.9%), Zr (>99.9%) and Al (99.99%). Metallic glass ribbons were synthesized by rapid quenching from the molten state using a single roller technique under Ar gas atmosphere. The density of the metallic glass was measured by gas pycnometer (Micro-metric, Accupyc 1330) with He gas (99.999%) at room temperature.

The neutron diffraction experiment was carried out by using the High Intensity Total scattering spectrometer HIT-II at the High Energy Accelerator Research Organization (KEK), Tsukuba, Japan. The measurement of the structure factor, $S(Q)$, over a wide Q range (from $Q = 5 \text{ nm}^{-1}$ up to 200 nm^{-1} in this case), leading to a high resolution in real-space, can be obtained by time-of-flight neutron diffraction using a spallation neutron source. The samples (about 2 g) were put into the vanadium cells (8 mm in inside diameter, 40 mm in height and 0.025 mm in wall thickness) for measuring the scattering intensity. The intensity was corrected for the cell, background, absorption and, multiple and incoherent scattering. The obtained coherent scattering intensity, $I(Q)$, was normalized with the scattering intensity from a vanadium rod as standard.

The X-ray diffraction measurement was carried out with photon energy of 113.68 keV at the high-energy X-ray diffraction beam line BL04B2 in the SPring-8, Mikazuki-cho, Hyogo, Japan. The sample was put into a flat container of 2 mm thickness with Kapton windows. After corrections for polarization, absorption and Compton scattering, the $I(Q)$ was converted into $S(Q)$, over a wide Q range (from $Q = 5 \text{ nm}^{-1}$ up to 200 nm^{-1}).

The $S(Q)$ of a multi-component system according to the Faber-Ziman definition⁴⁾ is derived as follows;

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Fig. 1 Structure factors, $S(Q)$, observed by neutron diffraction (solid line) and Reverse Monte Carlo (RMC) fits (dashed line) for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ (a) and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ (b) metallic glasses.

$$I(Q) = N\langle b \rangle^2 S(Q) + N(\langle b^2 \rangle - \langle b \rangle^2), \quad (1)$$

where $I(Q)$ is the coherent intensity, Q is the magnitude of the scattering vector, N is number of atoms and

$$\langle b^2 \rangle = \sum_i C_i b_i^2, \quad \langle b \rangle = \sum_i C_i b_i, \quad (2)$$

in which C_i and b_i are respectively the concentration and the coherent neutron scattering length (atomic scattering factor for X-ray) of element i . The structure factor,

$$S(Q) = \frac{1}{\langle b \rangle^2} \sum_{i,j=1}^n C_i C_j b_i b_j S_{ij}(Q), \quad (3)$$

where $S_{ij}(Q)$ is the partial structure factor (i and j each refer to either of the elements in the sample). $S_{ij}(Q)$ is related to the partial pair correlation function, $g_{ij}(r)$, for correlations between atoms of element i and of element j , by Fourier transformation,

$$Q(S_{ij}(Q) - 1) = 4\pi\rho_o \int_0^{Q_{\max}} r(g_{ij}(r) - 1)\sin(Qr)dr, \quad (4)$$

where ρ_o is the atomic number density of the sample.

Reverse Monte Carlo (RMC) modeling³⁾ based on the diffraction data is a simple variant of the standard Monte Carlo procedure, but a powerful method for deriving three-dimensional atomic arrangements of disordered materials. The atoms in a periodic model are moved at random and the moves are accepted or rejected in such a way as to minimize the difference between experiment and simulation, expressed as

$$\chi^2 = \sum_{i=1}^m [S_{\text{rmc}}(Q_i) - S_{\text{exp}}(Q_i)]^2 / \sigma^2(Q_i), \quad (5)$$

where $S_{\text{exp}}(Q)$ is the structure factor obtained experimentally,

and $S_{\text{rmc}}(Q)$ is the structure factor for the model. $\sigma(Q_i)$ is the experimental error for the i -th data point, $S(Q_i)$, determined from counting statistics.

In this study, we started with an initial configuration of about 5000 atoms, randomly distributed in a cube of side $L \sim 4.6$ nm. Normal periodic boundary conditions were used and L was chosen to be sufficiently large that there is little truncation effect in the Fourier transformation of eq. (4). The number of atoms and cube size were chosen so that the model had the experimental density. In this RMC modeling, various kinds of initial atomic configurations were used to confirm the reliability of the model structure. Moreover, the simulation was carried out by use of both neutron and X-ray diffraction data simultaneously. These procedures improve the confidence of the model structure. Accordingly, a small structural change observed in the study could be reconfirmed. The local atomic environment for the simulations was characterized by means of a Voronoi analysis.⁵⁾ The radii connecting an atom with its neighbours are bisected with planes and then the intersection of these planes forms the Voronoi polyhedron of the atom. The distribution of the Voronoi polyhedra is expressed in terms of Schläfli notation⁵⁾ whereby a polyhedron is denoted by a set of indices $n_3, n_4, n_5, n_6, n_7, n_8$, where n_i is the number of faces with i vertices. Note that the total coordination number of a specific atom is given by the sum of these indices, $\sum_i n_i$.

3. Results

Figure 1 shows the structure factors, $S(Q)$, observed by neutron and X-ray diffraction (solid lines) for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ (a) and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ (b) metallic glasses. The characteristic pre-peak is clearly observed at $Q \sim 17 \text{ nm}^{-1}$ in $S(Q)$ s for both $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. However, it

Fig. 2 Partial structure factors, $S_{ij}(Q)$, derived from the Reverse Monte Carlo (RMC) simulations for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ (a) and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ (b) metallic glasses.

can be found that the intensity of the pre-peak decreased by alloying Al to the Ni-Zr system and the amplitude of the oscillations above $Q \sim 60 \text{ nm}^{-1}$ for the Ni-Zr-Al metallic glass became smaller than for the Ni-Zr metallic glass. Moreover, there is no sign of a small angle scattering within the range of the experimental data ($Q_{\min} = 5 \text{ nm}^{-1}$), indicating that the ternary metallic glass is homogeneous.

In order to get information of the three-dimensional atomic arrangement for these metallic glasses, RMC modeling based on the neutron and X-ray diffraction data was employed. Figure 1 also indicates the $S(Q)$ derived from the RMC models (dashed lines), and excellent agreement with the experimental $S(Q)$ observed by neutron and X-ray diffraction. The pre-peak observed at $Q \sim 17 \text{ nm}^{-1}$ in the neutron $S(Q)$ for both $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses was explicitly reproduced by the RMC models. Partial structure factors, $S_{ij}(Q)$, derived from the RMC model for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses are shown in Fig. 2. The pre-peak in the total $S(Q)$ for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses was found to arise from the Ni-Ni partial structure factor, $S_{\text{NiNi}}(Q)$. But it can be seen that the intensity of the pre-peak for $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass becomes lower in comparison with that for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass due to addition of Al.

Fig. 3 Bond angle distributions for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ (solid line) and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ (dashed line) metallic glasses.

4. Discussion

When the pre-peak for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass is compared with that for $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass, we can easily recognize that the addition of Al to the Ni-Zr system causes the change of the atomic configuration, definitely, as shown in Fig. 1. Moreover, partial $S_{ij}(Q)$ obtained the RMC model apparently indicates that the structural change is attributed to the Ni-Ni correlation, since there is no drastic change in the partial $S_{ij}(Q)$ s of Ni-Zr and Zr-Zr pair correlations as seen in Fig. 2. The result teaches us that when Al is alloyed with the Ni-Zr system, Al atoms affect principally the Ni-Ni atomic arrangement. Figure 3 shows bond angle distributions for Ni-Ni-Ni, Ni-Ni-Zr, Zr-Ni-Zr, Ni-Zr-Zr, Ni-Zr-Ni and Zr-Zr-Zr derived from the RMC simulations for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. All of the bond angle distributions show two peaks, one in the region of $50\text{--}60^\circ$, and the other in the region of $100\text{--}120^\circ$. It can be seen that the Ni-Ni-Ni and Zr-Ni-Zr bond angles changed after adding Al to the Ni-Zr system. Therefore, the result also indicates that the Al atoms change the structure around a Ni atom in the Ni-Zr glass.

Voronoi analysis applied to investigate the polyhedral environment in the RMC model for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. Voronoi polyhedra around Ni, Zr and Al atoms were analyzed independently, as shown in Fig. 4. The fractions of the polyhedra around Zr atoms for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass are quite similar to those for $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass as shown in Fig. 4(a). The

Fig. 4 Fractions of Voronoi polyhedra around Zr (a), Ni (b) and Al (b) atoms for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses.

polyhedra around Zr atoms predominantly have total coordination numbers in the range of 12–14 atoms, typical of icosahedron-like geometry. The fractions of the polyhedra around Zr atoms are quite similar with each other for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. The same result has been reported in the previous paper,⁶⁾ which was discussed the structure of $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Cu}_{33.3}\text{Zr}_{66.7}$ metallic glasses. Therefore, we can easily recognize that the fundamental structure is constructed by the packing of Zr atoms with icosahedron-like environment for both $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. These polyhedra have various kinds of pentagonal symmetry planes, the fivefold symmetry of which is recognized to be important for stabilizing the glassy state. Thus far, the structural study has also suggested the existence of the icosahedral atomic configuration from the coordination number of 12.5 atoms around Zr in $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass.²⁾

The Voronoi polyhedra around Ni atoms differ evidently between $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses as shown in Fig. 4(b). For $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass the Voronoi polyhedra around Ni atoms are dominated by the trigonal prism capped with three half-octahedra (0 3 6 0 0 0), the Archimedean antiprism capped with two half-octahedra (0 2 8 0 0 0), the tetragonal dodecahedron (0 4 4 0 0 0), and analogous polyhedra. The coordination numbers for these dominant Voronoi polyhedra are respectively 9 atoms for (0 3 6 0 0 0) polyhedron, 8 atoms for (0 4 4 0 0 0) polyhedron and 10 atoms for (0 2 8 0 0 0) polyhedron, with an average of less than 10 atoms. These polyhedra correspond well with the structure of crystalline NiZr_2 .⁷⁾ Moreover, this figure clearly indicates that after adding Al to Ni-Zr system the polyhedra characterized by the trigonal prism drastically decrease in the $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass. It is worth noting that the fractions of the polyhedra with the coordination number more

Fig. 5 The proportion of icosahedron-like, prism-like polyhedra and the others around Zr, Ni and Al atoms for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses.

than 11 increase due to the addition of Al. Figure 4(c) shows the polyhedra around Al atoms for $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass. It can be recognized that the polyhedra of icosahedron-like geometry with more than 11 atoms is dominant in the structure around Al atoms. The Voronoi analysis of the atomic arrangement of the RMC model for $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glass definitely demonstrates an increase of the icosahedron-like polyhedra which have higher coordination number greater than 11 atoms.

Figure 5 shows the proportion of the icosahedron-like polyhedra {(0 0 12 0 0 0), (0 2 8 2 0 0), (0 2 8 1 0 0)}, prism-like polyhedra {(0 4 4 0 0 0), (0 3 6 0 0 0), (0 2 8 0 0 0)} and

the others in the RMC model for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses. The result allows us to clearly understand the increase of the icosahedron-like and the drastic decrease of the prism-like polyhedra after alloying Al with the Ni-Zr system. Especially, since the decrease of the prism-like polyhedra around Ni atoms $\{(0\ 4\ 4\ 0\ 0\ 0), (0\ 3\ 6\ 0\ 0\ 0), (0\ 2\ 8\ 0\ 0\ 0)\}$ is pronounced, it is likely that the stability of the Ni-Zr-Al metallic glasses is strongly affected by decreasing the polyhedra correspond well with the structure of crystalline NiZr_2 as well as by increasing the icosahedron-like polyhedra. The result is fairly supported by our previous result that icosahedron-like polyhedra dominate for the environment around Cu atoms for $\text{Cu}_{33.3}\text{Zr}_{66.7}$ metallic glass in comparison with for that around Ni atoms for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass.⁶⁾

5. Conclusions

The Voronoi analysis of the model structure obtained by RMC modeling based on the neutron and X-ray diffraction data for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ and $\text{Ni}_{25}\text{Zr}_{60}\text{Al}_{15}$ metallic glasses was investigated from the viewpoint of the stability of the glassy state. It informed us that the Zr atoms form a very similar fundamental structure with icosahedron-like geometry in both systems and such polyhedra with a total coordination number in excess of 12 play a role to form the glassy configuration. The polyhedra around Ni atoms are dominated by trigonal prism-like, Archimedian antiprism-like, and similar polyhedra, which correspond with the structure of

the crystalline NiZr_2 , with total coordination numbers less than 10 for $\text{Ni}_{33.3}\text{Zr}_{66.7}$ metallic glass. However, after alloying Al with the Ni-Zr system, the drastic decrease of the prism-like polyhedra around Ni atoms and the domination of the icosahedron-like polyhedra around Zr and Al atoms were observed. Therefore, these results indicate that icosahedron-like polyhedra play an important role for stabilizing the glass configuration.

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